

# Measurement of $\gamma$ -ray Emission Probabilities of $^{100}\text{Tc}$

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## Abstract

In order to precisely determine thermal neutron capture cross section of  $^{99}\text{Tc}$  nuclide,  $\gamma$ -ray emission probabilities of  $^{100}\text{Tc}$  were measured using a newly developed  $\beta$ - $\gamma$  coincidence system, which utilizes a plastic scintillator as a  $\beta$ -ray detector.

About 5 kBq of  $^{99}\text{Tc}$  samples were irradiated for 15 seconds, and then  $\beta$  and  $\gamma$  rays were measured for three minutes. The measurements were repeated 99 times.

From the obtained data, ratios of coincidence counting rates to  $\beta$  singles counting rates,  $n_c/n_\beta$ , was determined for 540 keV  $\gamma$  ray. To determine  $\gamma$ -ray emission probabilities, a method has been developed which uses a Monte Carlo simulation and takes into account  $\beta$  feedings to excited levels and detection threshold of  $\beta$  rays.

## 1 Introduction

The  $^{99}\text{Tc}$  is a long-lived ( $T_{1/2} = 2.1 \times 10^5$  y) fission product nuclide with a large fission yield. Therefore, it is suggested as a candidate for neutron transmutation, and many experimental works have been done in order to determine the thermal neutron capture cross sections of  $^{99}\text{Tc}$ , in particular by using an activation method[1].

In order to precisely determine the cross section in the activation method, precise values of  $\gamma$ -ray emission probabilities of the reaction product,  $^{100}\text{Tc}$ , are required. These values are, however, determined with large errors ( $7.0 \pm 1.2$  % for 540 keV  $\gamma$  ray [2]) mainly because of its short half-life ( $T_{1/2} = 15.5$  s).

To determine  $\gamma$ -ray emission probabilities of the  $^{100}\text{Tc}$  nuclide, a  $\beta$ - $\gamma$  coincidence measurement has been done by using a newly developed  $\beta$ - $\gamma$  coincidence system which utilizes a plastic scintillator as a  $\beta$ -ray detector[3].

## 2 The experiment and data analysis

The experiment has been carried out at the Research Reactor Institute of Kyoto University. About 10  $\mu\text{l}$  of ammonium hydroxide solution which contained 5 kBq of  $^{99}\text{Tc}$  as ammonium pertechnetate, was dried on an acrylic plate and irradiated by neutrons for 15 seconds using a pneumatic tube Pn-3 at the institute, and then  $\beta$ - and  $\gamma$ -rays emitted from the sample were measured.

Beta rays were measured using a plastic scintillation detector placed at 10 cm distance from the sample, whose thickness was 4 mm and which has a trapezoidal shape with an area of 33  $\text{cm}^2$ . For  $\gamma$ -ray detection, a HPGe detector was used, whose relative efficiency is 90 % of that of 3"  $\times$  3" NaI detector. The distance between the irradiated sample and the Ge detector was 3 cm. To follow the decay of the  $^{100}\text{Tc}$  nuclei, the obtained data were partitioned every 5 seconds. The irradiations and the measurements were repeated 99 times to improve the statistical accuracy. A total of  $2 \times 10^6$  counts were obtained for a peak area of the 540 keV  $\gamma$ -ray, which is emitted after  $\beta$  decay of the  $^{100}\text{Tc}$  nuclei.

In **figure 1**, an energy spectrum of  $\gamma$  rays observed by the Ge detector is shown, which is obtained by summing singles (projection) spectra over all runs. From the spectrum,  $\gamma$  rays emitted following the decay of  $^{100}\text{Tc}$ , namely 540, 591 and 1512 keV, are clearly identified. The figure also shows a

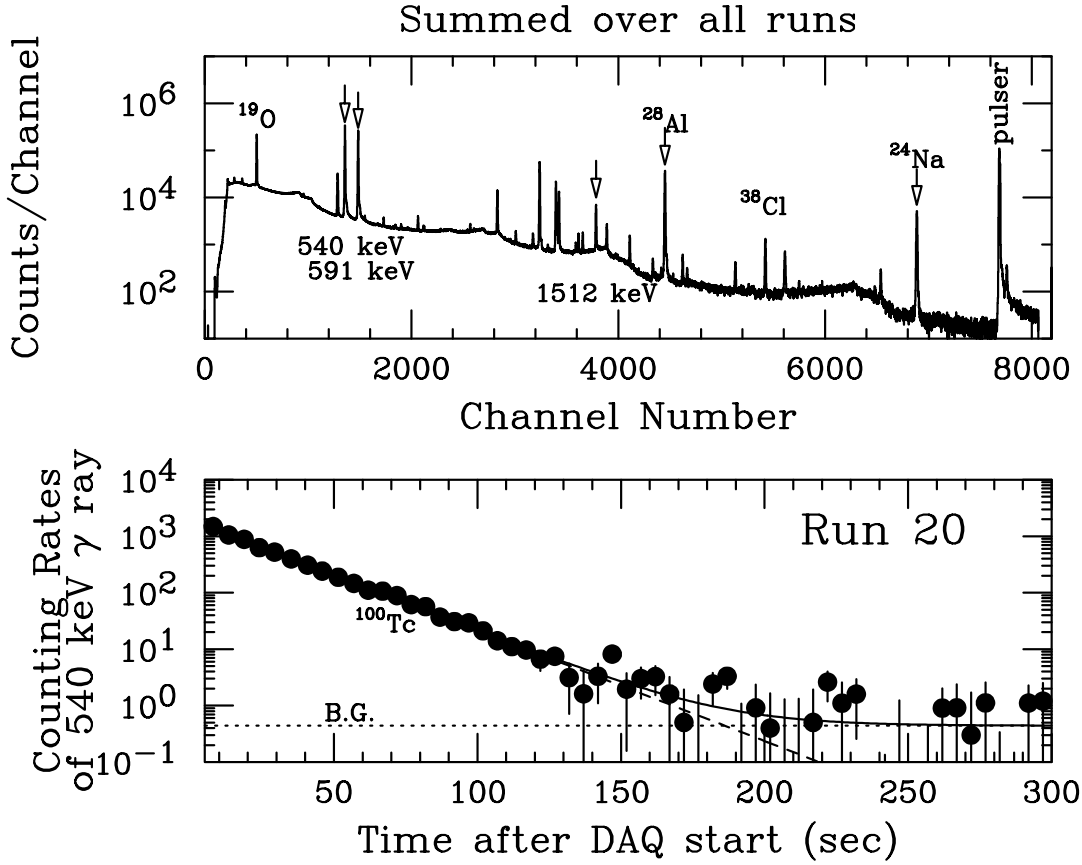


Figure 1: (upper) Singles spectrum of  $\gamma$  rays summed over all runs. (lower) Decay curve of 540 keV peak yields. The dashed line represents contribution from  $^{100}\text{Tc}$ , while dotted lines are that from backgrounds.

decay curve of peak-areal counts of the 540 keV  $\gamma$  ray in a typical run. The decay data for each run were fitted by a function with the following form:

$$Y(t) = a \exp\left(-\frac{\ln(2)}{\tau}t\right) + C, \quad (1)$$

where  $\tau$ ,  $a$  and  $C$  were varied as free parameters. The  $\tau$  corresponds to the half-life of  $^{100}\text{Tc}$ . By averaging the obtained values for  $\tau$  over all runs,  $\tau$  for  $^{100}\text{Tc}$  was determined to be  $15.27 \pm 0.02$  (s), thus the 540 keV  $\gamma$  rays were confirmed to be emitted following the decay of  $^{100}\text{Tc}$ .

Shown in **figure 2** is a  $\beta$ -ray energy spectrum summed over all runs. A bump around 50 channel bin was caused by noises, and therefore counts above 90 channel (inclusive) are summed as a  $\beta$  counts. **Figure 3** shows a decay curve for the  $\beta$ -ray counting rate in a typical run. To extract a contribution from  $^{100}\text{Tc}$  to the singles  $\beta$  counts, the decay curves were fitted by a function with the following form:

$$Y'(t) = \sum_{i=1}^3 a_i \exp\left(-\frac{\ln(2)}{\tau_i}t\right) + C', \quad (2)$$

In the fit,  $a_i$ s and  $C'$  were varied as free parameters while  $\tau_i$ s were fixed to half-lives of  $^{100}\text{Tc}$  (15.27 s),  $^{19}\text{O}$  (26.91 s [2]) and  $^{27}\text{Al}$  (2.248 min. [3]). The solid line in the figure shows a result of a fit, while contribution from the decay of  $^{100}\text{Tc}$  was shown by the dashed line. Coincidence counting rates of the  $\beta$  rays were determined in the same manner.

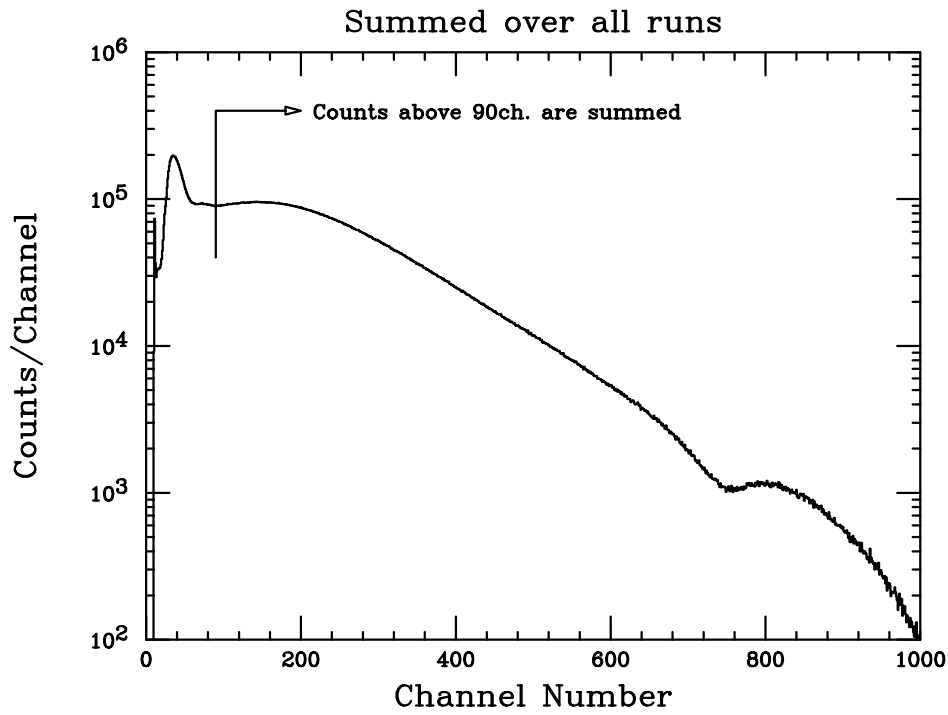


Figure 2: Singles spectrum of  $\beta$  rays summed over all runs.

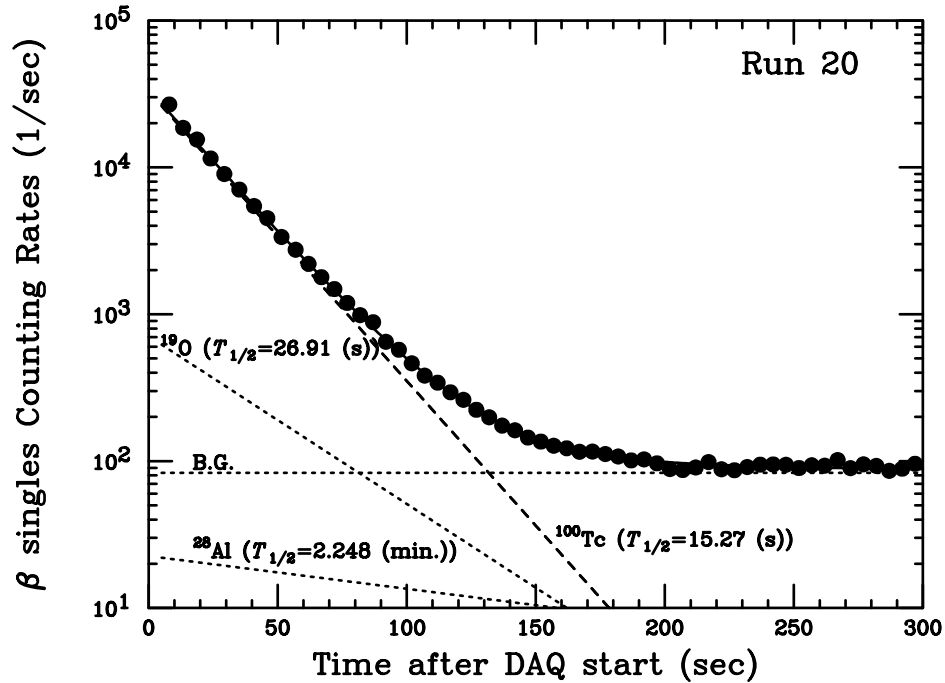


Figure 3: Decay curves of singles  $\beta$  ray counts above 90 ch. The solid line represents the result of the fit and dashed line contribution from  $^{100}\text{Tc}$ . The dotted lines express contributions from  $^{19}\text{O}$ ,  $^{28}\text{Al}$  and backgrounds.

From the result of the above mentioned fits, ratios of coincidence counting rates to  $\beta$  singles counting rates,  $n_c/n_\beta$ , were calculated for each time bin in each run. The final values of  $n_c/n_\beta$  were obtained after averaging over all time bins and all runs. For 540 keV  $\gamma$  ray,  $n_c(540\text{ keV})/n_\beta$  was determined to be  $0.00168 \pm 0.00001$ .

To deduce  $\gamma$ -ray emission probabilities from the obtained  $n_c/n_\beta$  values, a method has been developed using the Monte Carlo simulation code EGS4[4]. It simulates the decay of the  $^{100}\text{Tc}$  nuclei and the interaction between the detectors and the emitted particles ( $\beta$  and  $\gamma$  rays), and calculates  $n_c/n_\beta$  values as a function of  $\beta$  branching ratio to the ground state of  $^{100}\text{Ru}$ . Four excited states of  $^{100}\text{Ru}$  with large  $\beta$  branching ratios were taken into account in the simulation, namely the ground state, 540, 1131, 1362 and 2052 keV levels. With the five levels, more than 99.8 % of  $\beta$  decay feedings are exhausted[2]. Ratios between the feeding probabilities to the excited states were determined from the obtained singles  $\gamma$ -ray spectrum, after sum-coincidence corrections. By varying the  $\beta$  feeding probability to the ground state and searching for the value which reproduce the observed  $n_c/n_\beta$  values, the  $\gamma$ -ray emission probabilities can be determined. The analysis is now under way.

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